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Grant N00014-90-J-1193

TECHNICAL REPORT No. 8

Reply to "Coupled s- and d-Wave States in the Thorium-Doped Heavy-Fermion Superconductor  $UBe_{13}$ " by H. Pleiner and H. R. Brand

by

D. Sahu, A. Langner and Thomas F. George

Prepared for Publication

in

Physical Review B

Departments of Chemistry and Physics  
State University of New York at Buffalo  
Buffalo, New York 14260

March 1990

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REPORT DOCUMENTATION PAGE				Form Approved OMB No. 0704-0188	
1a. REPORT SECURITY CLASSIFICATION Unclassified			1b. RESTRICTIVE MARKINGS		
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2b. DECLASSIFICATION/DOWNGRADING SCHEDULE					
4. PERFORMING ORGANIZATION REPORT NUMBER(S) UBUFFALO/DC/90/TR-8			5. MONITORING ORGANIZATION REPORT NUMBER(S)		
6a. NAME OF PERFORMING ORGANIZATION Depts. Chemistry & Physics State University of New York		6b. OFFICE SYMBOL (If applicable)	7a. NAME OF MONITORING ORGANIZATION		
6c. ADDRESS (City, State, and ZIP Code) Fronczak Hall, Amherst Campus Buffalo, New York 14260			7b. ADDRESS (City, State, and ZIP Code) Chemistry Program 800 N. Quincy Street Arlington, Virginia 22217		
8a. NAME OF FUNDING/SPONSORING ORGANIZATION Office of Naval Research		8b. OFFICE SYMBOL (If applicable)	9. PROCUREMENT INSTRUMENT IDENTIFICATION NUMBER Grant N00014-90-J-1193		
8c. ADDRESS (City, State, and ZIP Code) Chemistry Program 800 N. Quincy Street Arlington, Virginia 22217			10. SOURCE OF FUNDING NUMBERS		
			PROGRAM ELEMENT NO.	PROJECT NO.	TASK NO.
					WORK UNIT ACCESSION NO.
11. TITLE (Include Security Classification) Reply to "Coupled s- and d-Wave States in the Thorium-Doped Heavy-Fermion Superconductor UBe <sub>13</sub> " by H. Pleiner and H. R. Brand					
12. PERSONAL AUTHOR(S) D. Sahu, A. Laugner and Thomas F. George					
13a. TYPE OF REPORT		13b. TIME COVERED FROM _____ TO _____		14. DATE OF REPORT (Year, Month, Day) March 1990	
15. PAGE COUNT 4					
16. SUPPLEMENTARY NOTATION Prepared for publication in Physical Review B					
17. COSATI CODES			18. SUBJECT TERMS (Continue on reverse if necessary and identify by block number)		
FIELD	GROUP	SUB-GROUP	SUPERCONDUCTOR UBe <sub>13</sub> COUPLED s- AND d-WAVE STATES		
			THORIUM-DOPED THEORETICAL ANALYSIS		
			HEAVY FERMION SYMMETRY CONSIDERATIONS		
19. ABSTRACT (Continue on reverse if necessary and identify by block number)  We discuss the points raised by H. Pleiner and H. R. Brand concerning our earlier work on UBe <sub>13</sub> and thorium-doped UBe <sub>13</sub> .					
20. DISTRIBUTION/AVAILABILITY OF ABSTRACT <input checked="" type="checkbox"/> UNCLASSIFIED/UNLIMITED <input checked="" type="checkbox"/> SAME AS RPT. <input type="checkbox"/> DTIC USERS			21. ABSTRACT SECURITY CLASSIFICATION Unclassified		
22a. NAME OF RESPONSIBLE INDIVIDUAL Dr. David L. Nelson			22b. TELEPHONE (Include Area Code) (202) 696-4410		22c. OFFICE SYMBOL

DD Form 1473, JUN 86

Previous editions are obsolete.

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Physical Review B, in press

Reply to "Coupled s- and d-Wave States in the Thorium-Doped

Heavy-Fermion Superconductor UBe<sub>13</sub>"

by H. Pleiner and H. R. Brand

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Abstract

We discuss the points raised by H. Pleiner and H. R. Brand concerning our earlier work on UBe<sub>13</sub> and thorium-doped UBe<sub>13</sub>.

PACS Nos.: 74.70.Tx, 74.20.De

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In two recent papers<sup>1,2</sup> we examined the question of lowest-order coupling among even-parity superconducting states utilizing general symmetry considerations as guiding principles. Our approach was motivated by the work of Kumar and Wölfle (KW),<sup>3</sup> whose conclusions we generalized<sup>1</sup> and examined more closely.<sup>2</sup> Our works<sup>1,2</sup> have been criticized by Pleiner and Brand<sup>4</sup> (PB) who argue that we misinterpreted KW. Our response spells out why the criticisms of PB appear to us to be without foundation.

PB claim<sup>4</sup> that KW followed a two-step procedure for minimizing their free energy in Ref. 3: first KW minimized their free energy with respect to  $\Delta_1$  (the order parameter for  $d_{x^2-y^2}$ -state) and concluded that this must be zero for stability reasons. This enabled KW, in a second step, to write down their free energy involving  $\Delta_2$  (the order parameter for the axial  $d_{3z^2-x^2-y^2}$ -state) and  $\Delta_0$  (the order parameter for the isotropic s-state). Since this two-step procedure is never mentioned by KW in writing Eq. 1 of Ref. 3, the question of misinterpreting KW on this point does not arise.

In addition, we would like to comment on the claim of PB that a non-zero  $\Delta_1$  can never be an equilibrium state. Our detailed numerical work<sup>2</sup> suggests otherwise. We find that a transition from an s-wave state (state IV of Ref. 2) to a mixed state with both d-wave components of the two-dimensional representation non-zero (state VII of Ref. 2) is clearly favored over that of a transition to a mixed state with  $\Delta_1 = 0$ ,  $\Delta_2 \neq 0$  and  $\Delta_0 \neq 0$  (state VI of Ref. 2). To account for the specific heat data, one clearly needs a non-zero  $\Delta_1$ , as has been realized by KW<sup>5</sup> and by us.<sup>6</sup>

We are well aware of the required symmetry breaking to achieve second-order transitions to the various superconducting states. This point has been well addressed in Ref. 7. However, as is also mentioned in Ref. 7, the

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symmetry breaking is to a lower symmetry subgroup of the cubic group. Our analysis incorporates this progression. It should be noted that PB do not substantiate their claims by any detailed calculations as we have done in Ref. 2; they merely state in general terms that in the symmetry-broken phase "a given state can have lower symmetry than the crystal symmetry." The question as to whether state VI of Ref. 2 is stable or unstable with respect to the introduction of a  $\Delta_1$  perturbation is left unanswered by PB. Furthermore, their earlier work in Ref. 8, being based on the unsubstantiated points of Ref. 4, seems to us to be questionable.

This research was supported by the Office of Naval Research and by an award from the New York State Institute on Superconductivity in conjunction with the New York State Energy Research and Development Authority.

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